# NOVEL ANALYTIC GRADIENT EVALUATION TECHNIQUES FOR OPTIMIZATION OF MICROWAVE STRUCTURES 

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## Introduction

ACCURATE design of modern and increasingly complex microwave devices and structures can often be achieved only through sophisticated optimization methods [1], [2]. It is generally accepted that efficient optimization methods are those exploiting not only the values of the cost function but its gradient and even its Hessian matrix as well. The gradient of the cost function is most of the time evaluated by repeated analyses at closely-spaced values of the independent variables; such a process can be prohibitively time consuming when the number of independent variables is large.
Recently, the Adjoint Network Method (ANM), which has been known for a long while in the area of circuit optimization, has been successfully used to analytically evaluate the gradient of cost functions for optimization of microwave filters and radiating structures [3], [4]. By its own nature, the ANM requires a network representation of the structure to be optimized (and its adjoint). The Mode-Matching Technique (MMT) was used to extract the network's matrix representation in form of, e.g., the admittance matrix [3] or the scattering matrix [4]. Starting from the scattering matrix of a network, including internal ports, a direct method of evaluating network sensitivities (gradient) was presented in [5].
Unfortunately more general numerical techniques such as the Method of Moments (MoM) [6] or the Finite Element Method (FEM) [7] do no directly yield a network representation in general.

Within these two methods, scattering problems are reduced to matrix equations of the form

$$
\begin{equation*}
[A][x]=[b] . \tag{1}
\end{equation*}
$$

Here, $[A]$ is a $M \times M$ matrix which depends on the independent variables and represents the structure to be optimized, $[b]$ is the excitation and $[x]$ is the response. For example, the vector $[x]$ contains the expansion coefficients in MoM or the nodal values in FEM.
With these features in mind, the paper addresses the following questions:

1. Is it possible to determine analytically the gradient of cost and response functions with respect to structural changes in the system without first deriving a network representation from $[x]$ and without finite differencing?
2. Is it possible to calculate group delays analytically from the matrix $[A]$ and excitation $[b]$ ?
It is shown that, as long as the partial derivatives of the matrix $[A]$ and the excitation $[b]$ are known analytically, all relevant sensitivities can also be determined analytically.

## Theory

We assume that the analysis of the structure to be optimized has been put in the form of equation (1). The optimization problem consists in determining the optimal values of the independent variables $a_{1}, a_{2}, \ldots, a_{n}$ such that the response function is equal to an optimal function $x^{o p t}$. Note that the independent variables appear explicitly in both the matrix $[A]$ as well as the excitation $[b]$
in general. We also assume that the optimization problem consists in minimizing a cost function $F$ which depends explicitly on the solution $[x]$.
The partial derivative of the cost function with respect to $a_{i}$ can be written in the form

$$
\begin{equation*}
\frac{\partial F}{\partial a_{i}}=\sum_{j=1}^{M} \frac{\partial F}{\partial x_{j}} \frac{\partial x_{j}}{\partial a_{i}} \tag{2}
\end{equation*}
$$

Since the partial derivatives $\frac{\partial F}{\partial x_{j}}$ are known, or can be straightforwardly determined, the problem reduces to calculating the partial derivatives of the solution $[x]$ with respect to $a_{i}$. To evaluate these, we take the partial derivative or equation (1) to get

$$
\begin{equation*}
\frac{\partial[A]}{\partial a_{i}}[x]+[A] \frac{\partial[x]}{\partial a_{i}}=\frac{\partial[b]}{\partial a_{i}} . \tag{3}
\end{equation*}
$$

Here, the derivative of a matrix is defined as a matrix whose elements are the derivatives of the original matrix. Rearranging this last equation and using equation (1), we get

$$
\begin{equation*}
\frac{\partial[x]}{\partial a_{i}}=[A]^{-1} \frac{\partial[b]}{\partial a_{i}}-[A]^{-1} \frac{\partial[A]}{\partial a_{i}}[A]^{-1}[b] . \tag{4}
\end{equation*}
$$

This equation shows that the partial derivatives are all known once the inverse $[A]^{-1}$ is known along with the partial derivatives of the entries of the matrix $[A]$ and $[b]$. If the partial derivatives of $[A]$ and $[b]$ are known analytically, which is the case of a large number of structures, the gradient is then also known analytically from a single inversion of the matrix $[A]$. In the more general case, the partial derivatives of $[A]$ and $[b]$ can be computed from finite differences and still determine the gradient of the cost function from a single inversion of the matrix $[A]$, or equivalently from a single analysis of the structure.
Situations may arise where the matrix $[A]$ itself involves inversion of other matrices. In such a case, the following expression for the partial derivatives of the inverse of a matrix in terms of those of the matrix itself is to be used

$$
\begin{equation*}
\frac{\partial[B]^{-1}}{\partial a_{i}}=-[B]^{-1} \frac{\partial[B]}{\partial a_{i}}[B]^{-1} \tag{5}
\end{equation*}
$$

We finally note that, if needed, the Hessian matrix can also be obtained from the present approach.

## Results

The present approach is first applied to a simple structure to document the accuracy of the computed partial derivatives. We consider a single E-plane stub in a rectangular waveguide. The width of the stub is $L$. Assuming that the fundamental $\mathrm{TE}_{10}$ is incident on the stub, we propose to compute the partial derivative $\frac{\partial S_{11}}{\partial L}$ as a function of frequency using the present approach and conventional finite differencing. The formulation of this simple scattering problem within the Coupled-Integral-Equation Technique (CIET) is given in [8]. Edge-conditioned basis functions are used to accelerate convergence of the numerical solution. Figures 1 and 2 show the real and imaginary parts of this partial derivative as obtained from the present analytic approach (solid line) and finite difference with an increment $\delta L=a / 100$. It is evident that both calculations agree well except for small deviations for frequencies between 11 and 13.8 GHz . However, when the increment is reduced to $\delta L=a / 1000$, the two results agree and can not be distinguished as Figures 3 and 4 show.


Fig. 1. Real part of $\frac{\partial S_{11}}{\partial L}$ as a function of frequency. Solid line: present method, dashed line: finite difference. $L=a=2 b=b_{1}=19.05 \mathrm{~mm}, \delta L=\mathrm{a} / 100$.

The next example consists in a 7 stub E-plane structure. We first evaluate the partial derivatives of the scattering parameters of the $\mathrm{TE}_{10}$ mode with respect to the inner dimensions of the filters (13 in all). Table I shows these quantities as obtained


Fig. 2. Imaginary part of $\frac{\partial S_{11}}{\partial L}$ as a function of frequency. Solid line: present method, dashed line: finite difference. $L=a=2 b=b_{1}=19.05 \mathrm{~mm}, \delta L=\mathrm{a} / 100$.


Fig. 3. Real part of $\frac{\partial S_{11}}{\partial L}$ as a function of frequency. Solid line: present method, dashed line: finite difference. $L=a=2 b=b_{1}=19.05 \mathrm{~mm}, \delta L=\mathrm{a} / 1000$.
from the present approach as well as finite differencing for two values of the increment. It is again evident that the results from the conventional finite differencing converge to the results obtained from the proposed technique when small enough values of the increments in the independent variables are used. This constitutes an additional validation of this novel approach.
The present approach can also be applied to computing group delays. It is easily shown that the group delay is given by [1]

$$
\begin{equation*}
\tau_{g}=-\operatorname{Im}\left[\frac{1}{S_{21}} \frac{\partial S_{21}}{\partial \omega}\right] \tag{6}
\end{equation*}
$$



Fig. 4. Imaginary part of $\frac{\partial S_{11}}{\partial L}$ as a function of frequency. Solid line: present method, dashed line: finite difference. $L=a=2 b=b_{1}=19.05 \mathrm{~mm}, \delta L=\mathrm{a} / 1000$.

The group delay of of the 7 -stub structure was


Fig. 5. Group delay of a 7-stub structure versus frequency. The analysis is performed using CIET with edge-conditioned basis functions [8].
computed using the present approach. Figure 5 shows that the results obtained from the present work (solid line) agree well with those from finite differencing the phase of $S_{12}$ (dashed line). These results were obtained from 2 edge-conditioned basis functions at each discontinuity.

Table I
Partial derivatives $\frac{\partial\left|S_{11}\right|}{\partial L_{i}}$

| i | This method | FD $^{1}$ | FD $^{2}$ |
| :---: | :---: | :---: | :---: |
| 1 | -126.4792 | -127.0087 | -126.4798 |
| 2 | 0.7323 | 0.8009 | 0.7324 |
| 3 | -126.9130 | -127.4637 | -126.9135 |
| 4 | 1.0517 | 1.1316 | 1.0518 |
| 5 | -127.1495 | -127.7106 | -127.1500 |
| 6 | 1.2017 | 1.2859 | 1.2018 |
| 7 | -127.2251 | -127.7893 | -127.2257 |
| 8 | 1.2017 | 1.2859 | 1.2018 |
| 9 | -127.1495 | -127.7106 | -127.1500 |
| 10 | 1.0517 | 1.1316 | 1.0518 |
| 11 | -126.9130 | -127.4637 | -126.9135 |
| 12 | 0.7323 | 0.8009 | 0.7324 |
| 13 | -126.4792 | -127.0087 | -126.4798 |

${ }^{1} \delta L_{i}=0.1 \mathrm{~mm}$
${ }^{2} \delta L_{i}=0.1 \mu \mathrm{~m}$

## Conclusions

A novel technique to evaluate gradients for optimization of microwave structures and devices was presented. The approach only requires that the problem be formulated in terms of a general nonhomogeneous matrix equation such as encountered in the Moment Method or the Finite Element Method. All partial derivatives are determined from a single analysis of the structure. Numerical results show excellent agreement between the present approach and the finite difference method when small enough increments are used.

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